

The Crystal and Molecular Structure of 4,4'-diamino-3,3'-dimethylbiphenyl (*o*-Tolidine)

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4,4'-Diamino-3,3'-dimethylbiphenyl, $C_{14}H_{16}N_2$, crystallizes with cell dimensions $a=7.44$, $b=23.70$, $c=6.47$ Å, space group $P2_12_12_1$ and $Z=4$. The structure has been determined from three-dimensional X-ray data and refined by the minimum-residual method, with isotropic temperature factors for individual atoms, to give a final R index of 11%. The two phenyl rings are not coplanar but are twisted around the linkage between them, C(1)–C(1'), so as to be mutually inclined at an angle of 41° ; in addition each phenyl ring is bent through an angle of about 3° away from the line C(1)–C(1'). The length of the bond C(1)–C(1') is 1.504 ± 0.013 Å.

Introduction

X-ray analyses of crystals of biphenyl (Hargreaves & Rizvi, 1962; Robertson, 1961*a, b*; Trotter, 1961) have established that the biphenyl molecule is planar in the solid state, although electron diffraction studies (Brockway & Karle, 1944; Bastiansen, 1949; Almenningen & Bastiansen, 1958) suggest that the two phenyl rings are mutually inclined at an angle of about 42° in the vapour phase. Substituents introduced in the 2,2' positions produce overcrowding which must result in a non-planar molecule (see *e.g.* Smare, 1948; Fowweather & Hargreaves, 1950; Fowweather, 1952). It is not clear, however, what will be the effect of introducing substituents in the 3,3' and 4,4' positions. A number of these derivatives are in course of examination by X-ray methods and we now present details of the stereochemistry of one of them, *viz.* 4,4'-diamino-3,3'-dimethylbiphenyl (*o*-tolidine).

Experimental

X-ray rotation and Weissenberg photographs about the three principal axes of the orthorhombic crystals show that the space group is $P2_12_12_1$ and that four molecules are contained in a unit cell with the following dimensions and estimated limits of error; $a=7.44 \pm 0.01$, $b=23.70 \pm 0.05$, $c=6.47 \pm 0.02$ Å. These results are in agreement with preliminary X-ray studies by Clark & Pickett (1931), Hargreaves (1940), Quader (1958) and Lasheen (1963).

Reflexions $hk0$, $hk1$, $hk2$, $hk3$ and $hk4$, recorded in equi-inclination Weissenberg photographs using Cu $K\alpha$ radiation, provided the data from which the atomic positions were finally determined. The intensities of the reflexions were measured by the multiple film technique and visual comparison with crystal reflected calibration

spots of known relative exposures. No corrections were made for absorption, which was small in the crystal specimen used for intensity measurements ($\mu=6.7$ cm^{-1} ; dimensions of cross-section perpendicular to the axis of oscillation ~ 0.18 mm).

Determination and refinement of the structure

Optical transforms and Patterson maps were used to obtain the approximate orientations of the phenyl rings and the position of the centre of the molecule was determined by the molecular location technique of Bhuiya & Stanley (1964). The [001] projection agreed well with that deduced previously (Sullivan, 1960) by direct methods using Sayre's relationship.

Refinement began with an agreement index, R , of 0.44 for the $hk0$ reflexions and 0.42 for the $0kl$ reflexions. Three cycles of Fourier refinement using ρ_o maps reduced $R(hk0)$ to 0.35. Because of overlap of the projected atoms refinement was continued by the minimum residual method of Bhuiya & Stanley (1963); six cycles of refinement reduced $R(hk0)$ to 0.12. $R(0kl)$ was similarly reduced to 0.19 by three cycles of Fourier refinement and to 0.15 by five cycles of minimum residual refinement.

At this stage ($\rho_o - \rho_c$) maps indicated that the structure was basically correct and that there was no marked thermal anisotropy. Minimum residual refinement was then applied to all the 839 observed hkl reflexions, excepting for 6 reflexions which appeared to suffer from extinction effects. The calculated structure factors were based on the atomic scattering factors of Forsyth & Wells (1959) and isotropic temperature factors B . Four cycles of refinement reduced $R(hkl)$ from 0.17 to 0.13. Hydrogen atoms were then included in the calculations and after five more cycles of refinement, when the parameter shifts were less than one-quarter of the estimated standard deviations, a final R index of 0.11 was obtained.

The final positional and thermal parameters are given in Table 1. Estimated standard deviations of the posi-

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tional coordinates were deduced with the use of Cruickshank's (1960) formula; the z coordinates required a correction for the fact that experimental data had not been collected to the edge of the limiting sphere in this direction. The standard deviations obtained thus were

	$\sigma(x)=\sigma(y)$	$\sigma(z)$
Carbon	0.008 Å	0.017 Å
Nitrogen	0.007	0.013

The observed and calculated structure amplitudes, and the calculated phase angles α , are presented in Table 2.

Description of the structure

Fig. 1 gives the dimensions of the *o*-tolidine molecule deduced from the positional coordinates in Table 1. The estimated standard deviations of the bond lengths and bond angles are given within brackets; the former were calculated with the use of Ahmed & Cruickshank's (1953) formula and the latter by Darlow's (1960) formula.

The equations of the best planes passing through the atoms in phenyl ring I [C(1), C(2), ... C(6)] and phenyl ring II [C(1'), C(2'), ... C(6')], respectively, are

Phenyl ring I

$$0.9128x + 0.2403y - 0.3302z - 3.2001 = 0$$

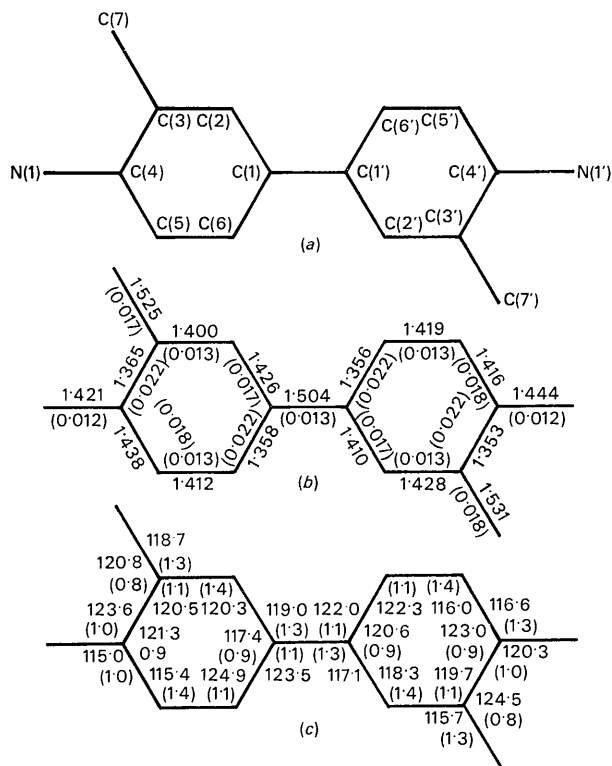


Fig. 1. Dimensions of the *o*-tolidine molecule. (a) Numbering of the atoms. (b) Bond lengths with standard deviations (Å). (c) Bond angles with standard deviations (°).

Table 1. Final positional coordinates and isotropic temperature factors

	x/a	y/b	z/c	B
C(1)	0.4119	0.1115	0.1097	2.60 Å ²
C(2)	0.4880	0.0705	0.2452	2.86
C(3)	0.5158	0.0151	0.1772	2.92
C(4)	0.4633	-0.0010	-0.0162	3.10
C(5)	0.3900	0.0390	-0.1605	3.40
C(6)	0.3647	0.0941	-0.0827	3.20
C(7)	0.5968	-0.0275	0.3271	3.58
N(1)	0.4847	-0.0566	-0.0944	3.98
C(1')	0.3775	0.1701	0.1897	2.75
C(2')	0.4060	0.2152	0.0515	2.97
C(3')	0.3646	0.2710	0.1197	2.99
C(4')	0.2962	0.2791	0.3111	3.16
C(5')	0.2688	0.2346	0.4536	3.50
C(6')	0.3095	0.1796	0.3809	3.28
C(7')	0.3989	0.3181	-0.0374	3.52
N(1')	0.2535	0.3353	0.3820	3.98
H(1)	0.331	0.120	-0.179	4.80
H(2)	0.360	0.023	-0.290	4.80
H(3)	0.421	-0.068	-0.232	4.80
H(4)	0.546	-0.076	0.027	4.80
H(5)	0.569	-0.031	0.464	4.80
H(6)	0.568	-0.063	0.277	4.80
H(7)	0.650	-0.014	0.465	4.80
H(8)	0.520	0.077	0.404	4.80
H(1')	0.298	0.151	0.490	4.80
H(2')	0.202	0.234	0.595	4.80
H(3')	0.177	0.335	0.509	4.80
H(4')	0.313	0.351	0.259	4.80
H(5')	0.303	0.321	-0.121	4.80
H(6')	0.384	0.356	-0.058	4.80
H(7')	0.459	0.345	0.079	4.80
H(8')	0.433	0.210	-0.108	4.80

Phenyl ring II

$$0.9300x + 0.1054y + 0.3522z - 3.4694 = 0$$

Deviations of atoms from these two planes are shown in Table 3 and Fig. 2.

Phenyl rings I and II are both planar within the accuracy of our measurements, but the molecule as a whole is non-planar since the two rings are mutually inclined at an angle, ϕ , of 41° . We may describe the relative positions of these rings by starting with a transplanar configuration and then bringing the rings to their observed mutual positions by introducing the three following operations: (i) rotation of one ring with respect to the other about the axis C(1)-C(1') through an angle ϕ_1 , (ii) rotation of ring I through an angle ϕ_2 about an axis which is in the plane of ring I and passes through C(1) in a direction at right angles to the line C(1)-C(1'), (iii) rotation of ring II through an angle ϕ_3 about an axis which is in the plane of ring II and passes through C(1') in a direction at right angles to the line C(1)-C(1'). The angles ϕ_2 and ϕ_3 are quite small (2.7° and 3.3° respectively) and because of this the angle between the rings, $\phi=41^\circ$, differs from ϕ_1 by only a few minutes. Fig. 2 shows that the angles ϕ_2 and ϕ_3 represent bends in the length of the molecule at C(1) and C(1'); it also shows, through plots of the atomic positions determined, that these bends are real

Table 2 (cont.)

Table with 3 columns of data (h k l |0|f| |0|f| |alpha|) repeated three times. Each column contains numerical values, including integers, decimals, and percentages, representing crystallographic data.

Table 3. Displacements from the best planes through the atoms in phenyl rings I and II respectively

Ring I		Ring II	
C(1)	0.001 Å	C(1')	0.003 Å
C(2)	-0.004	C(2')	-0.002
C(3)	0.015	C(3')	0.006
C(4)	-0.021	C(4')	-0.011
C(5)	0.017	C(5')	0.012
C(6)	-0.008	C(6')	-0.009
C(7)	-0.001	C(7')	0.006
N(1)	-0.028	N(1')	-0.004
C(1')	-0.070	C(1)	-0.087
C(4')	-0.262	C(4)	-0.299
N(1')	-0.381	N(1)	-0.468

and are not, in the main, to be attributed to uncertainties in the atomic positions.

It is interesting to compare the molecular configuration of *o*-tolidine with that of biphenyl, which is planar in the crystalline phase (see *Introduction*). Planarity requires an unusually close approach of the hydrogen atoms attached to the two pairs of atoms C(2), C(6') and C(6), C(2') and it would appear that, in biphenyl, crystal forces and the resonance energy of a planar structure are sufficient to overcome the effects of steric hindrance; the observed C(1)–C(1') distance of 1.497 ± 0.003 Å indicates that the phenyl rings are linked by a C(sp²)–C(sp²) single bond and suggests that the planarity of the molecule may be largely determined by crystal forces.

In *o*-tolidine the C(1)–C(1') distance is virtually the same, *viz.* 1.504 ± 0.013 Å, but it is likely that steric effects are enhanced by the amino and methyl groups in the 4,4' and 3,3' positions respectively and that the resulting stresses are reduced by the rotations φ_1 , φ_2 and φ_3 ; again crystal forces may play an important part in determining the configuration adopted. The hydrogen atoms attached to C(2) and C(6') are separated by a distance of 2.3 Å and those attached to C(6) and C(2') by 2.4 Å. These distances suggest there

is no serious strain due to steric interaction between the hydrogen atoms, but little significance can be attached to the value observed because of their very low precision – a very crude treatment gives an e.s.d. of 0.3 Å.

Non-coplanar phenyl rings have been reported in two similar biphenyl derivatives, *viz.* 4,4'-diamino-3,3'-dichlorobiphenyl [Hasan (1962) – the structure has been refined and a paper is now in preparation] and 4,4'-dinitrobiphenyl (Boonstra, 1963).

There are no significant differences (Cruickshank, 1949) between the lengths of equivalent bonds in the two halves of the molecule – although there appear to be some significant differences in bond angles – but there are marked differences between the bond lengths within each aromatic ring. Assuming that the two aromatic rings are identical, and taking mean values, we find that there are two short bond lengths [C(1)–C(6), 1.357 Å and C(3)–C(4), 1.359 Å] and four long ones [C(1)–C(2), 1.418 Å; C(2)–C(3), 1.414 Å; C(4)–C(5), 1.427 Å; and C(5)–C(6), 1.415 Å]. The differences between the long and short bond lengths appear to be significant, although some uncertainty exists because no libration corrections have been applied to the bond lengths. In biphenyl (Robertson, 1961*a,b*) the variations in corresponding bond lengths show much smaller variation, ranging from 1.344 to 1.385 Å when uncorrected for libration and from 1.378 to 1.398 when corrected for libration.

The carbon atoms of the methyl groups and the nitrogen atoms of the amino groups form bonds of the usual lengths with the carbon atoms to which they are attached in the aromatic rings and, within experimental error, lie in the plane of the aromatic ring.

Fig. 3 shows the structure viewed along the direction [001] and indicates all distances less than 4.1 Å between atoms in neighbouring molecules; the intermolecular distances are quite normal, the shortest being 3.44 Å.

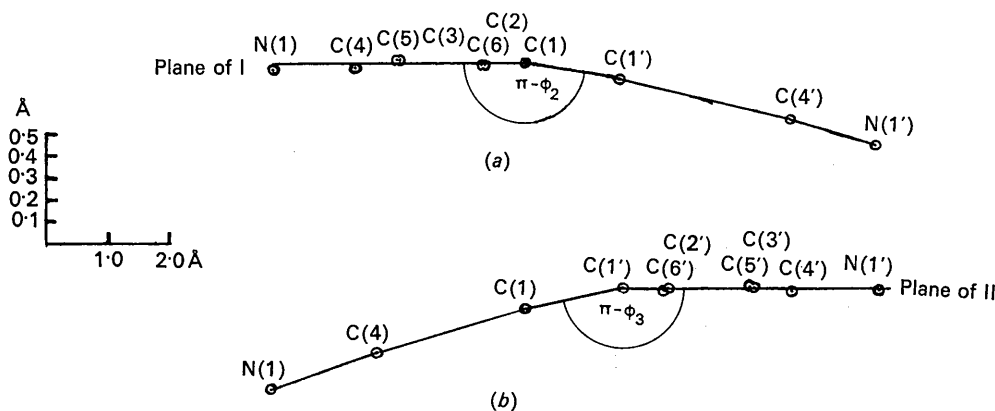


Fig. 2. Displacements from the best plane through atoms in (a) phenyl ring I, (b) phenyl ring II. In both (a) and (b) the best plane is viewed end-on and in a direction at right angles to the length of the molecule. The best planes are shown by horizontal lines and the displacements of individual atoms are measured by vertical displacements from this line; for clarity, vertical displacements are magnified by plotting them on a scale approximately three times as large as the horizontal scale.

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A New Class of Compound $M_x^+ A_x^{3+} Ti_{2-x} O_4$ ($0.60 < x < 0.80$) Typified by $Rb_x Mn_x Ti_{2-x} O_4$

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$Rb_x Mn_x Ti_{2-x} O_4$ is a non-stoichiometric compound with orthorhombic symmetry, space group *Imm2*, and at the composition $x=0.75$ has the unit-cell dimensions $a=3.934$, $b=15.918$, $c=2.933$ Å. Its structure contains (Ti, Mn)-oxygen octahedra sharing edges in puckered sheets normal to (010) similar to those in lepidocrocite, $FeO(OH)$, except that the correspondence between sheets has been altered to give Rb^+ an irregular eightfold coordination. A number of isomorphous compounds have been prepared and the field of occurrence so far examined, while restricted to the alkali metals Rb^+ and Cs^+ , includes a wide range of trivalent and divalent metals substituting for Ti^{4+} in the octahedral positions. The isomorphism of the member $Cs_x Mg_{x/2} Ti_{2-x/2} O_4$ is confirmed by a least-squares refinement of powder diffractometer data.

Introduction

Oxycompounds containing alkali, trivalent and quadrivalent metals offer a useful means of studying stereochemical relations in solid state chemistry. We recently showed that a number of compounds $NaA^{3+}B^{4+}O_4$ isomorphous with calcium ferrite could be synthesized, but only when *A* and *B* were metal ions of spherical symmetry averaged over two virtually identical, but crystallographically different octahedral positions (Reid, Wadsley & Sienko, 1967). Other classes of structure with different stoichiometries were generated when the trivalent ions A^{3+} ordered into positions of special symmetry. For example $NaTi_2Al_5O_{12}$ (Mumme & Wadsley, 1967) and $Na_4Mn_4Ti_5O_{18}$ (Mumme, 1968) formed instead of 'NaAlTiO₄' and 'NaMnTiO₄', where Al was present in tetrahedral as well as octahedral sites in the one, and Mn^{3+} octahedral and square pyramidal in the other. One other factor limiting the formation of calcium ferrite structures was the size of

the alkali metal ion, the available positions being too large for Li^+ and too small for K^+ , Rb^+ or Cs^+ .

When a stable compound of a new kind emerges from such experimental studies, the composition is often established more rapidly by determining its crystal structure than by isolating enough of it for chemical analysis; thereafter its confirmation by synthesis as a single phase is usually straightforward. However the possibility of a lattice compound being non-stoichiometric is a variable to be considered in a structure analysis, which provides *ab initio* an ideal composition based on point symmetry.

In the present study two classes of oxide phases of variable composition containing Rb^+ or Cs^+ , Ti^{4+} , and a number of tri- or divalent metals were recognized. The composition limits of both types were determined experimentally by treating each multicomponent system as a quasi-binary. Fig. 1 shows the $Rb-Ti-Mn^{3+}-O$ system reduced to the join $RbMnO_2-TiO_2$. One phase had the hollandite structure (Byström & Byström,