

The Crystal and Molecular Structure of 4,4'-diamino-3,3'-dimethylbiphenyl (*o*-Tolidine)

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(Received 25 October 1967)

4,4'-Diamino-3,3'-dimethylbiphenyl, $C_{14}H_{16}N_2$, crystallizes with cell dimensions $a=7.44$, $b=23.70$, $c=6.47$ Å, space group $P2_12_12_1$ and $Z=4$. The structure has been determined from three-dimensional X-ray data and refined by the minimum-residual method, with isotropic temperature factors for individual atoms, to give a final R index of 11%. The two phenyl rings are not coplanar but are twisted around the linkage between them, C(1)–C(1'), so as to be mutually inclined at an angle of 41°; in addition each phenyl ring is bent through an angle of about 3° away from the line C(1)–C(1'). The length of the bond C(1)–C(1') is 1.504 ± 0.013 Å.

Introduction

X-ray analyses of crystals of biphenyl (Hargreaves & Rizvi, 1962; Robertson, 1961*a,b*; Trotter, 1961) have established that the biphenyl molecule is planar in the solid state, although electron diffraction studies (Brockway & Karle, 1944; Bastiansen, 1949; Almenningen & Bastiansen, 1958) suggest that the two phenyl rings are mutually inclined at an angle of about 42° in the vapour phase. Substituents introduced in the 2,2' positions produce overcrowding which must result in a non-planar molecule (see e.g. Smare, 1948; Fowweather & Hargreaves, 1950; Fowweather, 1952). It is not clear, however, what will be the effect of introducing substituents in the 3,3' and 4,4' positions. A number of these derivatives are in course of examination by X-ray methods and we now present details of the stereochemistry of one of them, *viz.* 4,4'-diamino-3,3'-dimethylbiphenyl (*o*-tolidine).

Experimental

X-ray rotation and Weissenberg photographs about the three principal axes of the orthorhombic crystals show that the space group is $P2_12_12_1$ and that four molecules are contained in a unit cell with the following dimensions and estimated limits of error; $a=7.44 \pm 0.01$, $b=23.70 \pm 0.05$, $c=6.47 \pm 0.02$ Å. These results are in agreement with preliminary X-ray studies by Clark & Pickett (1931), Hargreaves (1940), Quader (1958) and Lasheen (1963).

Reflexions $hk0$, $hk1$, $hk2$, $hk3$ and $hk4$, recorded in equi-inclination Weissenberg photographs using Cu $K\alpha$ radiation, provided the data from which the atomic positions were finally determined. The intensities of the reflexions were measured by the multiple film technique and visual comparison with crystal reflected calibration

spots of known relative exposures. No corrections were made for absorption, which was small in the crystal specimen used for intensity measurements ($\mu=6.7$ cm⁻¹; dimensions of cross-section perpendicular to the axis of oscillation ~ 0.18 mm).

Determination and refinement of the structure

Optical transforms and Patterson maps were used to obtain the approximate orientations of the phenyl rings and the position of the centre of the molecule was determined by the molecular location technique of Bhuiya & Stanley (1964). The [001] projection agreed well with that deduced previously (Sullivan, 1960) by direct methods using Sayre's relationship.

Refinement began with an agreement index, R , of 0.44 for the $hk0$ reflexions and 0.42 for the $0kl$ reflexions. Three cycles of Fourier refinement using ϱ_o maps reduced $R(hk0)$ to 0.35. Because of overlap of the projected atoms refinement was continued by the minimum residual method of Bhuiya & Stanley (1963); six cycles of refinement reduced $R(hk0)$ to 0.12. $R(0kl)$ was similarly reduced to 0.19 by three cycles of Fourier refinement and to 0.15 by five cycles of minimum residual refinement.

At this stage $(\varrho_o - \varrho_c)$ maps indicated that the structure was basically correct and that there was no marked thermal anisotropy. Minimum residual refinement was then applied to all the 839 observed hkl reflexions, excepting for 6 reflexions which appeared to suffer from extinction effects. The calculated structure factors were based on the atomic scattering factors of Forsyth & Wells (1959) and isotropic temperature factors B . Four cycles of refinement reduced $R(hkl)$ from 0.17 to 0.13. Hydrogen atoms were then included in the calculations and after five more cycles of refinement, when the parameter shifts were less than one-quarter of the estimated standard deviations, a final R index of 0.11 was obtained.

The final positional and thermal parameters are given in Table 1. Estimated standard deviations of the posi-

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tional coordinates were deduced with the use of Cruickshank's (1960) formula; the z coordinates required a correction for the fact that experimental data had not been collected to the edge of the limiting sphere in this direction. The standard deviations obtained thus were

	$\sigma(x) = \sigma(y)$	$\sigma(z)$
Carbon	0.008 Å	0.017 Å
Nitrogen	0.007	0.013

The observed and calculated structure amplitudes, and the calculated phase angles α , are presented in Table 2.

Description of the structure

Fig. 1 gives the dimensions of the *o*-tolidine molecule deduced from the positional coordinates in Table 1. The estimated standard deviations of the bond lengths and bond angles are given within brackets; the former were calculated with the use of Ahmed & Cruickshank's (1953) formula and the latter by Darlow's (1960) formula.

The equations of the best planes passing through the atoms in phenyl ring I [C(1), C(2), ..., C(6)] and phenyl ring II [C(1'), C(2'), ..., C(6')], respectively, are

Phenyl ring I

$$0.9128x + 0.2403y - 0.3302z - 3.2001 = 0$$

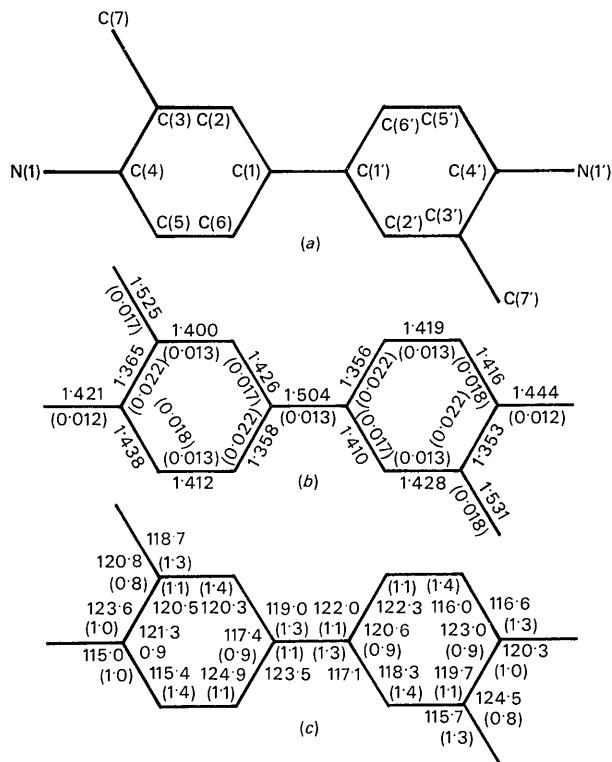


Fig. 1. Dimensions of the *o*-tolidine molecule. (a) Numbering of the atoms. (b) Bond lengths with standard deviations (Å). (c) Bond angles with standard deviations (°).

Table 1. Final positional coordinates and isotropic temperature factors

	x/a	y/b	z/c	B
C(1)	0.4119	0.1115	0.1097	2.60 Å ²
C(2)	0.4880	0.0705	0.2452	2.86
C(3)	0.5158	0.0151	0.1772	2.92
C(4)	0.4633	-0.0010	-0.0162	3.10
C(5)	0.3900	0.0390	-0.1605	3.40
C(6)	0.3647	0.0941	-0.0827	3.20
C(7)	0.5968	-0.0275	0.3271	3.58
N(1)	0.4847	-0.0566	-0.0944	3.98
C(1')	0.3775	0.1701	0.1897	2.75
C(2')	0.4060	0.2152	0.0515	2.97
C(3')	0.3646	0.2710	0.1197	2.99
C(4')	0.2962	0.2791	0.3111	3.16
C(5')	0.2688	0.2346	0.4536	3.50
C(6')	0.3095	0.1796	0.3809	3.28
C(7')	0.3989	0.3181	-0.0374	3.52
N(1')	0.2535	0.3353	0.3820	3.98
H(1)	0.331	0.120	-0.179	4.80
H(2)	0.360	0.023	-0.290	4.80
H(3)	0.421	-0.068	-0.232	4.80
H(4)	0.546	-0.076	0.027	4.80
H(5)	0.569	-0.031	0.464	4.80
H(6)	0.568	-0.063	0.277	4.80
H(7)	0.650	-0.014	0.465	4.80
H(8)	0.520	0.077	0.404	4.80
H(1')	0.298	0.151	0.490	4.80
H(2')	0.202	0.234	0.595	4.80
H(3')	0.177	0.335	0.509	4.80
H(4')	0.313	0.351	0.259	4.80
H(5')	0.303	0.321	-0.121	4.80
H(6')	0.384	0.356	-0.058	4.80
H(7')	0.459	0.345	0.079	4.80
H(8')	0.433	0.210	-0.108	4.80

Phenyl ring II

$$0.9300x + 0.1054y + 0.3522z - 3.4694 = 0.$$

Deviations of atoms from these two planes are shown in Table 3 and Fig. 2.

Phenyl rings I and II are both planar within the accuracy of our measurements, but the molecule as a whole is non-planar since the two rings are mutually inclined at an angle, φ , of 41°. We may describe the relative positions of these rings by starting with a transplanar configuration and then bringing the rings to their observed mutual positions by introducing the three following operations: (i) rotation of one ring with respect to the other about the axis C(1)-C(1') through an angle φ_1 , (ii) rotation of ring I through an angle φ_2 about an axis which is in the plane of ring I and passes through C(1) in a direction at right angles to the line C(1)-C(1'), (iii) rotation of ring II through an angle φ_3 about an axis which is in the plane of ring II and passes through C(1') in a direction at right angles to the line C(1)-C(1'). The angles φ_2 and φ_3 are quite small (2.7° and 3.3° respectively) and because of this the angle between the rings, $\varphi=41^\circ$, differs from φ_1 by only a few minutes. Fig. 2 shows that the angles φ_2 and φ_3 represent bends in the length of the molecule at C(1) and C(1'); it also shows, through plots of the atomic positions determined, that these bends are real

Table 2. Observed and calculated structure amplitudes and calculated phase angles (°)

<i>h</i>	<i>k</i>	<i>l</i>	<i>10/<i>f</i>₁</i>	<i>10/<i>f</i>₂</i>	α	<i>h</i>	<i>k</i>	<i>l</i>	<i>10/<i>f</i>₁</i>	<i>10/<i>f</i>₂</i>	α	<i>h</i>	<i>k</i>	<i>l</i>	<i>10/<i>f</i>₁</i>	<i>10/<i>f</i>₂</i>	α	<i>h</i>	<i>k</i>	<i>l</i>	<i>10/<i>f</i>₁</i>	<i>10/<i>f</i>₂</i>	α	
0	2	0	47	22	180.0	0	4	0	535	640	0.0	0	6	0	261	278	0.0	0	8	0	281	264	180.0	
0	12	0	259	248	0.0	0	14	0	107	95	180.0	0	16	0	110	103	180.0	0	18	0	317	314	0.0	
0	20	0	198	180	0.0	0	22	0	74	78	0.0	0	24	0	50	46	180.0	0	28	0	56	52	0.0	
*1	1	0	108	216	-90.0	*1	2	0	267	368	-190.0	1	3	0	31	50	0.0	1	4	0	41	74	-90.0	
1	5	0	278	295	-90.0	1	6	0	105	180	-90.0	1	7	0	137	144	-190.0	1	8	0	169	127	-90.0	
1	9	0	111	101	-90.0	1	11	0	102	106	-90.0	1	12	0	214	204	-190.0	1	13	0	134	128	-90.0	
1	14	0	125	108	-90.0	1	16	0	71	47	-90.0	1	17	0	63	50	-90.0	1	18	0	50	30	-90.0	
1	20	0	110	103	-90.0	1	21	0	04	82	-90.0	1	22	0	66	64	-90.0	1	27	0	42	11	-90.0	
1	29	0	68	50	-90.0	1	30	0	86	86	-90.0	*2	0	0	244	320	-90.0	*2	1	0	926	1231	0.0	
*2	2	0	880	1144	0.0	2	4	0	100	107	0.0	2	5	0	170	101	0.0	2	6	0	80	83	180.0	
2	7	0	222	228	0.0	2	8	0	04	65	180.0	2	9	0	20	21	180.0	2	10	0	193	139	0.0	
2	11	0	222	202	0.0	2	12	0	08	88	180.0	2	13	0	232	203	0.0	2	14	0	187	172	0.0	
2	16	0	199	185	0.0	2	17	0	109	102	180.0	2	19	0	107	87	0.0	2	20	0	83	68	0.0	
2	21	0	154	162	180.0	2	22	0	71	79	180.0	2	25	0	36	34	180.0	2	27	0	74	54	180.0	
2	28	0	48	43	0.0	3	1	0	301	294	90.0	3	2	0	204	195	90.0	3	3	0	162	170	-90.0	
3	4	0	45	43	-90.0	3	5	0	119	112	-90.0	3	6	0	94	85	-90.0	3	8	0	143	126	-90.0	
3	10	0	225	211	-90.0	3	12	0	88	87	90.0	3	14	0	127	109	-90.0	3	16	0	163	162	-90.0	
3	18	0	119	107	-90.0	3	19	0	273	271	90.0	3	20	0	56	52	-90.0	3	21	0	102	95	-90.0	
3	23	0	110	104	-90.0	3	24	0	53	48	90.0	3	25	0	50	41	90.0	3	26	0	53	43	90.0	
3	28	0	45	43	-90.0	3	29	0	39	33	90.0	4	1	0	63	65	180.0	4	2	0	148	145	0.0	
4	3	0	123	135	0.0	4	4	0	110	102	0.0	4	9	0	29	12	0.0	4	10	0	125	117	180.0	
4	11	0	119	96	0.0	4	12	0	145	147	0.0	4	14	0	228	227	0.0	4	16	0	166	167	0.0	
4	18	0	42	44	0.0	4	19	0	42	39	180.0	4	20	0	45	40	0.0	4	22	0	88	68	180.0	
4	23	0	107	92	180.0	4	25	0	45	33	180.0	4	27	0	39	26	180.0	5	1	0	39	33	-90.0	
5	2	0	80	60	-90.0	5	3	0	113	110	-90.0	5	5	0	50	42	33	90.0	6	0	0	32	32	-90.0
5	7	0	102	100	-90.0	5	8	0	116	121	-90.0	5	9	0	182	173	90.0	5	10	0	91	93	0.0	
5	11	0	134	132	-90.0	5	12	0	98	86	-90.0	5	13	0	98	78	90.0	5	14	0	86	68	-90.0	
5	15	0	94	95	90.0	5	16	0	150	98	100	5	17	0	98	102	90.0	5	18	0	193	193	-90.0	
5	20	0	110	98	90.0	5	21	0	60	53	90.0	5	22	0	66	54	-90.0	5	23	0	42	39	-90.0	
5	24	0	42	35	90.0	5	26	0	29	21	90.0	6	0	0	80	68	180.0	6	6	0	42	53	180.0	
6	2	0	130	125	0.0	6	3	0	102	102	180.0	6	5	0	103	53	45	180.0	6	11	0	39	54	180.0
6	7	0	77	75	180.0	6	8	0	53	43	180.0	6	9	0	53	38	180.0	6	16	0	91	83	0.0	
6	12	0	77	71	0.0	6	13	0	94	81	180.0	6	14	0	107	108	0.0	7	5	0	39	44	90.0	
6	17	0	32	41	0.0	7	6	0	32	33	-90.0	7	7	0	46	49	-90.0	7	10	0	86	62	90.0	
7	2	0	45	46	90.0	7	8	0	53	52	-90.0	7	9	0	83	102	-90.0	7	15	0	63	59	90.0	
7	11	0	94	88	90.0	7	12	0	56	56	-90.0	7	14	0	56	55	-90.0	8	0	0	119	144	0.0	
7	16	0	86	74	-90.0	7	17	0	68	48	-90.0	7	19	0	48	31	-90.0	8	9	0	48	45	180.0	
8	1	0	88	105	180.0	8	3	0	29	41	180.0	8	7	0	53	58	180.0	8	15	0	36	9	0.0	
8	10	0	53	55	180.0	8	11	0	98	79	180.0	8	12	0	53	51	0.0	8	15	0	36	9	0.0	
9	2	0	29	26	-90.0	9	3	0	36	47	90.0	9	2	0	31	242	217	180.0	9	3	1	203	192	90.0
5	1	27	26	90	0.0	5	6	1	114	115	0.0	5	7	1	89	66	-90.0	5	8	1	22	3	180.0	
5	15	1	38	37	-90.0	5	10	1	553	439	0.0	5	11	0	566	472	90.0	5	12	1	83	77	180.0	
5	15	1	134	127	-90.0	5	14	1	165	165	0.0	5	15	1	105	86	90.0	5	16	1	35	34	180.0	
5	17	1	51	45	-90.0	5	18	1	225	224	180.0	5	20	1	118	136	180.0	5	22	1	65	64	180.0	
5	23	1	51	47	-90.0	1	9	1	504	633	-90.0	1	1	1	787	870	-129.2	1	2	1	184	182	-160.1	
1	3	1	47	46	137.6	1	4	1	211	208	-140.8	1	5	1	141	126	-143.0	1	6	1	255	244	-23.1	
1	7	1	205	158	-59.9	1	8	1	174	107	-65.9	1	9	1	295	231	-177.5	1	10	1	130	118	-60.7	
1	11	1	115	97	91.4	1	12	1	172	132	-107.8	1	13	1	95	77	-132.4	1	14	1	56	53	-94.0	
1	15	1	123	97	-155.4	1	16	1	161	45	-137.5	1	17	1	120	107	-108.2	1	18	1	216	202	-111.1	
1	19	1	269	277	-140.2	1	20	1	199	211	128.6	1	21	1	54	49	-85.5	1	22	1	119	125	-50.7	
1	23	1	39	46	174.2	1	24	1	49	53	107.5	2	1	1	741	909	-90.0	2	1	1	528	576	-118.2	
2	2	1	305	366	-175.3	2	3	1	572	575	-145.0	2	4	1	177	170	-156.5	2	5	1	165	169	-114.2	
2	6	1	180	140	-174.1	2	7	1	146	136	30.0	2	8	1	46	54	-12.5	2	9	1	201	179	-147.4	
2	10	1	183	147	-82.9	2	11	1	255	245	10.4	2	12	1	143	133	16.5	2	13	1	192	172	-99.4	
2	14	1	99	95	136.0	2	15	1	81	78	33.6	2	16	1	94	96	118.6	2	17	1	136	142	-18.5	
2	18	1	92	93	-126.2	2	19	1	79	81	61.9	2	20	1	134	152	-177.2	2	21	1	57	50	-137.9	
3	4	1	61	56	90.0	3	5	1	206	216	171.0	3	6	1	306	400	-97.5	3	7	1	153	150	-135.6	
3	4	1	114	97	-139.4	3	5	1	272	277	-170.7	3	6	1	91	93	56.1	3	7	1	163	164	-126.1	
3	8	1	251	281	-8.0	3	10	1	165	154	-152.5	3	11	1	37	42	-138.3	3	12	1	51	49	-157.9	
3	13	1	73	67	165.8	3	14	1	59	45	-76.1	3	15	1	74	73	-166.7	3	16	1	31	33	-73.5	
3	17	1	153	161	-171.2	3	18	1	131	131	-30.4	3	19	1	293	264	-197.4	3	20	1	57	59	-173.4	
3	21	1	40	34	-147.7	7	9	1	77	84	-35.3	7</												

Table 2 (cont.)

h	k	l	$ O _{E_1}$	$ O _{E_2}$	α	h	k	l	$ O _{E_1}$	$ O _{E_2}$	α	h	k	l	$ O _{E_1}$	$ O _{E_2}$	α	h	k	l	$ O _{E_1}$	$ O _{E_2}$	α
5	2	2	11	88	-173.2	5	3	2	146	142	13.9	5	4	2	111	213	61.2	5	5	2	152	155	63.9
5	6	2	72	61	-11.0	5	7	2	86	86	41.7	5	9	2	62	65	74.5	5	10	2	84	92	162.5
5	11	2	25	47	-171.5	5	12	2	56	63	-7.5	5	13	2	53	52	125.5	5	14	2	43	49	-13.1
5	15	2	25	13	-137.2	5	16	2	112	106	88.2	5	17	2	56	42	47.2	5	18	2	68	71	-2.2
5	19	2	81	85	-169.0	5	20	2	19	21	121.1	5	21	2	52	47	19.7	5	23	2	23	32	33.7
6	0	2	16	98	180.0	6	1	2	35	54	-122.4	6	2	2	50	74	9.0	6	3	2	153	168	-11.5
6	4	2	60	65	-101.2	6	5	2	53	44	118.2	6	6	2	50	59	60.5	6	7	2	31	43	43.2
6	6	2	65	61	-104.5	6	9	2	64	81	1.1	6	10	2	53	87	107.3	6	11	2	31	43	-102.9
6	12	2	65	61	-104.5	6	13	2	54	57	8.0	6	14	2	41	45	156.4	6	17	2	76	82	-87.7
6	18	2	61	62	-105.5	6	19	2	70	51	-14.9	6	20	2	50	50	37.0	7	0	2	82	80	0.0
7	1	2	74	78	33.8	7	2	2	55	85	8.1	7	3	2	36	87	175.0	7	4	2	25	31	-6.6
7	5	2	61	82	61.3	7	6	2	84	79	85.5	7	7	2	77	84	50.4	7	8	2	24	22	-38.4
7	9	2	33	37	-41.6	7	10	2	61	51	31.2	7	11	2	53	44	130.9	7	12	2	31	51	-109.7
7	14	2	28	20	77.4	7	17	2	52	46	70.1	7	18	2	53	52	46.6	8	0	2	30	23	180.0
7	1	2	26	42	-108.3	8	2	2	50	53	-140.8	8	3	2	26	40	-67.8	8	4	2	23	49	86.4
8	5	2	32	41	9.0	8	6	2	32	46	144.9	8	7	2	68	77	81.6	8	8	2	66	70	76.4
8	9	2	71	80	-62.7	8	10	2	18	23	-75.8	8	11	2	71	85	71.9	8	12	2	16	39	-142.4
8	14	2	12	23	-101.5	9	0	2	35	38	0.0	9	1	2	34	11	53.2	9	2	2	52	71	-21.3
9	3	2	19	37	72.4	9	4	2	10	25	148.5	9	2	3	34	42	180.0	9	3	3	205	213	-99.0
0	4	3	127	126	0.0	0	5	3	233	229	-90.0	0	6	3	136	138	180.0	0	7	3	17	13	-90.0
0	8	3	61	44	0.0	0	9	3	274	205	-90.0	0	10	3	166	142	180.0	0	11	3	110	97	-90.0
0	12	3	16	26	180.0	0	13	3	141	148	-90.0	0	14	3	183	0.0	0	15	3	19	18	-90.0	
0	16	3	70	88	0.0	0	17	3	40	25	-90.0	0	18	3	20	33	180.0	0	21	3	24	21	-90.0
0	22	3	19	28	180.0	0	23	3	19	29	90.0	0	24	3	37	33	180.0	0	25	3	40	39	90.0
0	26	3	66	51	0.0	0	27	3	48	42	-90.0	0	28	3	42	39	180.0	1	0	3	138	190	90.0
1	1	3	105	218	-59.4	1	2	3	293	330	50.7	1	3	3	187	187	123.1	1	4	3	112	116	-160.2
1	5	3	44	35	171.4	1	7	3	38	18	-42.6	1	8	3	219	219	-82.9	1	9	3	159	165	-149.8
1	10	3	60	75	-101.1	1	11	3	49	34	65.2	1	12	3	133	144.3	1	13	3	97	93	-175.2	
1	14	3	73	65	45.3	1	15	3	91	72	152.3	1	16	3	149	140.2	1	17	3	53	43	47.8	
1	18	3	29	20	-120.0	1	19	3	82	89	-10.0	1	20	3	28	28	49.5	1	21	3	20	29	-175.0
1	22	3	24	31	156.1	1	25	3	39	36	121.1	1	26	3	43	40	-58.4	1	27	3	41	30	-136.7
1	28	3	18	10	-52.7	2	2	3	150	112	90.0	2	3	108	79	98.1	2	2	3	120	144	115.6	
2	3	3	137	144	-153.4	2	4	3	39	31	-146.1	2	5	3	110	105	8.5	2	6	3	97	75	-57.3
2	7	3	135	135	-82.7	2	8	3	73	82	163.1	2	9	3	48	53	-54.9	2	10	3	150	153	-144.9
2	11	3	78	80	-97.8	2	12	3	61	58	-113.0	2	13	3	146	139	-141.1	2	14	3	136	144	-69.0
2	15	3	94	82	-61.2	2	16	3	114	102	-55.0	2	17	3	29	34	-1.3	2	18	3	95	90	-9.2
2	19	3	76	78	-88.0	2	20	3	66	66	-152.3	2	21	3	68	58	-166.8	2	22	3	53	43	-101.5
2	23	3	18	12	174.3	2	24	3	20	27	157.9	2	25	3	37	37	-141.4	2	26	3	50	39	134.0
2	27	3	56	52	2.7	3	0	3	215	190	90.0	3	1	3	64	65	100.5	3	2	3	57	57	51.1
3	3	0	86	65	16.9	3	4	3	134	124	82.2	3	5	3	78	92	-99.5	3	6	3	102	110	-152.6
3	7	3	142	134	95.4	3	8	3	111	112	-38.5	3	9	3	45	39	-87.0	3	10	3	185	180	-152.7
3	11	3	139	148	-89.5	3	12	3	121	120	0.4	3	13	3	153	146	99.8	3	14	3	131	136	150.9
3	15	3	20	33	-63.0	3	16	3	125	123	121.0	3	17	3	54	48	-120.9	3	18	3	61	54	90.3
3	19	3	72	62	-150.1	3	20	3	53	52	-128.4	3	21	3	53	57	-107.3	3	23	3	62	60	174.2
3	24	3	43	39	-82.6	3	25	3	46	46	154.1	3	26	3	27	24	83.0	4	0	3	53	39	90.0
4	1	3	96	88	0.85	4	2	3	30	28	-23.6	4	3	2	64	65	110.9	4	4	3	76	78	-138.2
4	5	3	57	58	-174.9	4	6	3	54	57	-138.9	4	7	3	73	66	80.6	4	8	3	93	93	-44.4
4	9	3	100	94	-30.0	4	10	3	70	60	-140.0	4	11	3	28	36	-86.1	4	12	3	87	87	-51.5
4	13	3	35	49	163.1	4	14	3	86	83	-122.6	4	15	3	45	45	-123.4	4	16	3	95	91	-38.4
4	17	3	121	124	-106.7	4	18	3	19	19	-23.6	4	19	3	30	32	2.2	4	20	3	92	87	-29.8
4	21	3	56	54	-114.0	4	22	3	65	65	-92.3	4	23	3	56	56	-102.5	4	24	3	39	46	-162.3
5	0	3	53	49	-161.2	5	1	3	53	53	-12.5	5	2	3	113	114	-12.0	5	3	3	65	65	178.5
5	8	3	67	64	-161.2	5	3	3	157	138	1.0	5	4	3	116	124	-174.6	5	12	3	86	88	-120.8
5	12	3	97	109	-31.5	5	14	3	121	120	77.4	5	15	3	53	50	-14.5	5	16	3	27	27	174.6
5	17	3	89	80	144.8	5	18	3	81	70	78.1	5	19	3	24	49	-149.0	5	22	3	56	56	-76.1
5	23	3	37	35	178.4	6	0	3	160	167	-90.0	6	1	3	161	167	-104.0	6	2	3	70	70	-40.9
6	3	3	69	71	145.5	6	4	3	80	87	-76.4	6	5	3	47	52	65.2	6	6	3	76	84	-86.0
6	7	3	29	29	-144.5	6	8	3	21	35	-117.5	6	9	3	26	34	-8.3	6	10	3	38	35	-82.2
6	7	3	56	66	-26.2	6	10	3	50	58	-106.1	6	11	3	11	11	-90.0	6	12	3	21	21	0.0
6	14	4	76	91	-136.9	6	14	4	100	120	43.0	6	15	4	22	22	136.3	6	16	4	99	86	-89.5
6	18	4	55	47	133.6	6	19	4	48	39	93.2	6	20	4	108	108	-175.7	6	21	4	27	27	-7.1
6	22	4	27	47	50.7	6	13	4	236	260	144.8	6	14	4	42	42	-138.2	6	15	4	26	24	-147.7
6	16	4	52	56	-45.6	6	17	4	95	94	-177.4	6	18	4	42	42	-135.6	6	19	4	24	24	-144.6
6	24																						

Table 3. Displacements from the best planes through the atoms in phenyl rings I and II respectively

	Ring I	Ring II	
C(1)	0.001 Å	C(1')	0.003 Å
C(2)	-0.004	C(2')	-0.002
C(3)	0.015	C(3')	0.006
C(4)	-0.021	C(4')	-0.011
C(5)	0.017	C(5')	0.012
C(6)	-0.008	C(6')	-0.009
C(7)	-0.001	C(7')	0.006
N(1)	-0.028	N(1')	-0.004
C(1')	-0.070	C(1)	-0.087
C(4')	-0.262	C(4)	-0.299
N(1')	-0.381	N(1)	-0.468

and are not, in the main, to be attributed to uncertainties in the atomic positions.

It is interesting to compare the molecular configuration of *o*-tolidine with that of biphenyl, which is planar in the crystalline phase (see *Introduction*). Planarity requires an unusually close approach of the hydrogen atoms attached to the two pairs of atoms C(2), C(6') and C(6), C(2') and it would appear that, in biphenyl, crystal forces and the resonance energy of a planar structure are sufficient to overcome the effects of steric hindrance; the observed C(1)-C(1') distance of 1.497 ± 0.003 Å indicates that the phenyl rings are linked by a $C(sp^2)-C(sp^2)$ single bond and suggests that the planarity of the molecule may be largely determined by crystal forces.

In *o*-tolidine the C(1)-C(1') distance is virtually the same, *viz.* 1.504 ± 0.013 Å, but it is likely that steric effects are enhanced by the amino and methyl groups in the 4,4' and 3,3' positions respectively and that the resulting stresses are reduced by the rotations φ_1 , φ_2 and φ_3 ; again crystal forces may play an important part in determining the configuration adopted. The hydrogen atoms attached to C(2) and C(6') are separated by a distance of 2.3 Å and those attached to C(6) and C(2') by 2.4 Å. These distances suggest there

is no serious strain due to steric interaction between the hydrogen atoms, but little significance can be attached to the value observed because of their very low precision – a very crude treatment gives an e.s.d. of 0.3 Å.

Non-coplanar phenyl rings have been reported in two similar biphenyl derivatives, *viz.* 4,4'-diamino-3,3'-dichlorobiphenyl [Hasan (1962)] – the structure has been refined and a paper is now in preparation] and 4,4'-dinitrobiphenyl (Boonstra, 1963).

There are no significant differences (Cruickshank, 1949) between the lengths of equivalent bonds in the two halves of the molecule – although there appear to be some significant differences in bond angles – but there are marked differences between the bond lengths within each aromatic ring. Assuming that the two aromatic rings are identical, and taking mean values, we find that there are two short bond lengths [C(1)-C(6), 1.357 Å and C(3)-C(4), 1.359 Å] and four long ones [C(1)-C(2), 1.418 Å; C(2)-C(3), 1.414 Å; C(4)-C(5), 1.427 Å; and C(5)-C(6), 1.415 Å]. The differences between the long and short bond lengths appear to be significant, although some uncertainty exists because no libration corrections have been applied to the bond lengths. In biphenyl (Robertson, 1961*a,b*) the variations in corresponding bond lengths show much smaller variation, ranging from 1.344 to 1.385 Å when uncorrected for libration and from 1.378 to 1.398 when corrected for libration.

The carbon atoms of the methyl groups and the nitrogen atoms of the amino groups form bonds of the usual lengths with the carbon atoms to which they are attached in the aromatic rings and, within experimental error, lie in the plane of the aromatic ring.

Fig. 3 shows the structure viewed along the direction [001] and indicates all distances less than 4.1 Å between atoms in neighbouring molecules; the intermolecular distances are quite normal, the shortest being 3.44 Å.

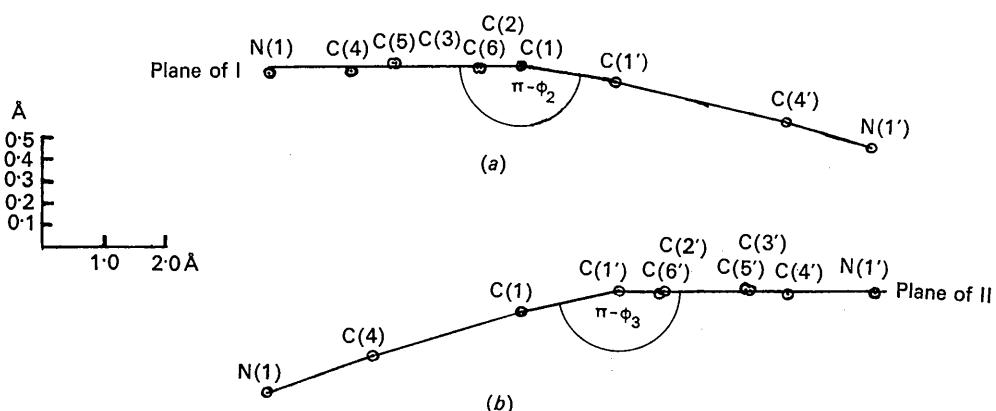


Fig. 2. Displacements from the best plane through atoms in (a) phenyl ring I, (b) phenyl ring II. In both (a) and (b) the best plane is viewed end-on and in a direction at right angles to the length of the molecule. The best planes are shown by horizontal lines and the displacements of individual atoms are measured by vertical displacements from this line; for clarity, vertical displacements are magnified by plotting them on a scale approximately three times as large as the horizontal scale.

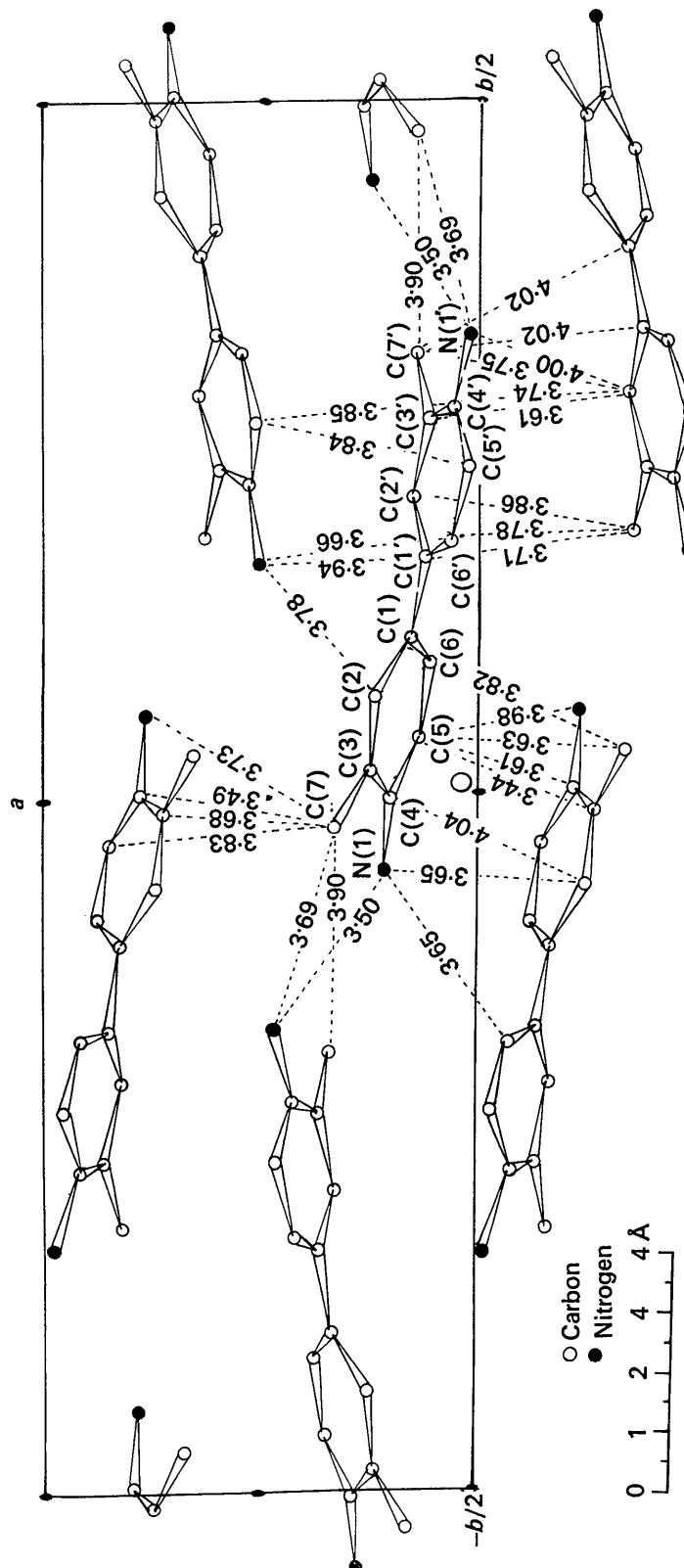


Fig. 3. The structure of *o*-toluidine viewed along [001], showing all intermolecular distances (excluding hydrogen atoms) of less than 4.1 Å.

One of us (S.A.C.) wishes to express thanks to the University of Rajshahi for study leave and to the British Council for a scholarship.

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A New Class of Compound $M_x^{+}A_x^{3+}Ti_{2-x}O_4$ ($0.60 < x < 0.80$) Typified by $Rb_xMn_xTi_{2-x}O_4$

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(Received 20 October 1967)

$Rb_xMn_xTi_{2-x}O_4$ is a non-stoichiometric compound with orthorhombic symmetry, space group *Imm2*, and at the composition $x=0.75$ has the unit-cell dimensions $a=3.934$, $b=15.918$, $c=2.933$ Å. Its structure contains (Ti, Mn)-oxygen octahedra sharing edges in puckered sheets normal to (010) similar to those in lepidocrocite, $FeO(OH)$, except that the correspondence between sheets has been altered to give Rb^+ an irregular eightfold coordination. A number of isomorphous compounds have been prepared and the field of occurrence so far examined, while restricted to the alkali metals Rb^+ and Cs^+ , includes a wide range of trivalent and divalent metals substituting for Ti^{4+} in the octahedral positions. The isomorphism of the member $Cs_xMg_{x/2}Ti_{2-x/2}O_4$ is confirmed by a least-squares refinement of powder diffractometer data.

Introduction

Oxycompounds containing alkali, trivalent and quadrivalent metals offer a useful means of studying stereochemical relations in solid state chemistry. We recently showed that a number of compounds $NaA^{3+}B^{4+}O_4$ isomorphous with calcium ferrite could be synthesized, but only when A and B were metal ions of spherical symmetry averaged over two virtually identical, but crystallographically different octahedral positions (Reid, Wadsley & Sienko, 1967). Other classes of structure with different stoichiometries were generated when the trivalent ions A^{3+} ordered into positions of special symmetry. For example $NaTi_2Al_5O_{12}$ (Mumme & Wadsley, 1967) and $Na_4Mn_4Ti_5O_{18}$ (Mumme, 1968) formed instead of ' $NaAlTiO_4$ ' and ' $NaMnTiO_4$ ', where Al was present in tetrahedral as well as octahedral sites in the one, and Mn^{3+} octahedral and square pyramidal in the other. One other factor limiting the formation of calcium ferrite structures was the size of

the alkali metal ion, the available positions being too large for Li^+ and too small for K^+ , Rb^+ or Cs^+ .

When a stable compound of a new kind emerges from such experimental studies, the composition is often established more rapidly by determining its crystal structure than by isolating enough of it for chemical analysis; thereafter its confirmation by synthesis as a single phase is usually straightforward. However the possibility of a lattice compound being non-stoichiometric is a variable to be considered in a structure analysis, which provides *ab initio* an ideal composition based on point symmetry.

In the present study two classes of oxide phases of variable composition containing Rb^+ or Cs^+ , Ti^{4+} , and a number of tri- or divalent metals were recognized. The composition limits of both types were determined experimentally by treating each multicomponent system as a quasi-binary. Fig. 1 shows the $Rb-Ti-Mn^{3+}-O$ system reduced to the join $RbMnO_2-TiO_2$. One phase had the hollandite structure (Byström & Byström,